

Topics

Nonlinear Equations

Analysis of PDEs

Elliptic PDEs & Boundary-Value ODEs

Parabolic PDEs & Initial Value ODEs

Nonlinear Equations

CHEN 6603

Nonlinear Eqns. - Overview



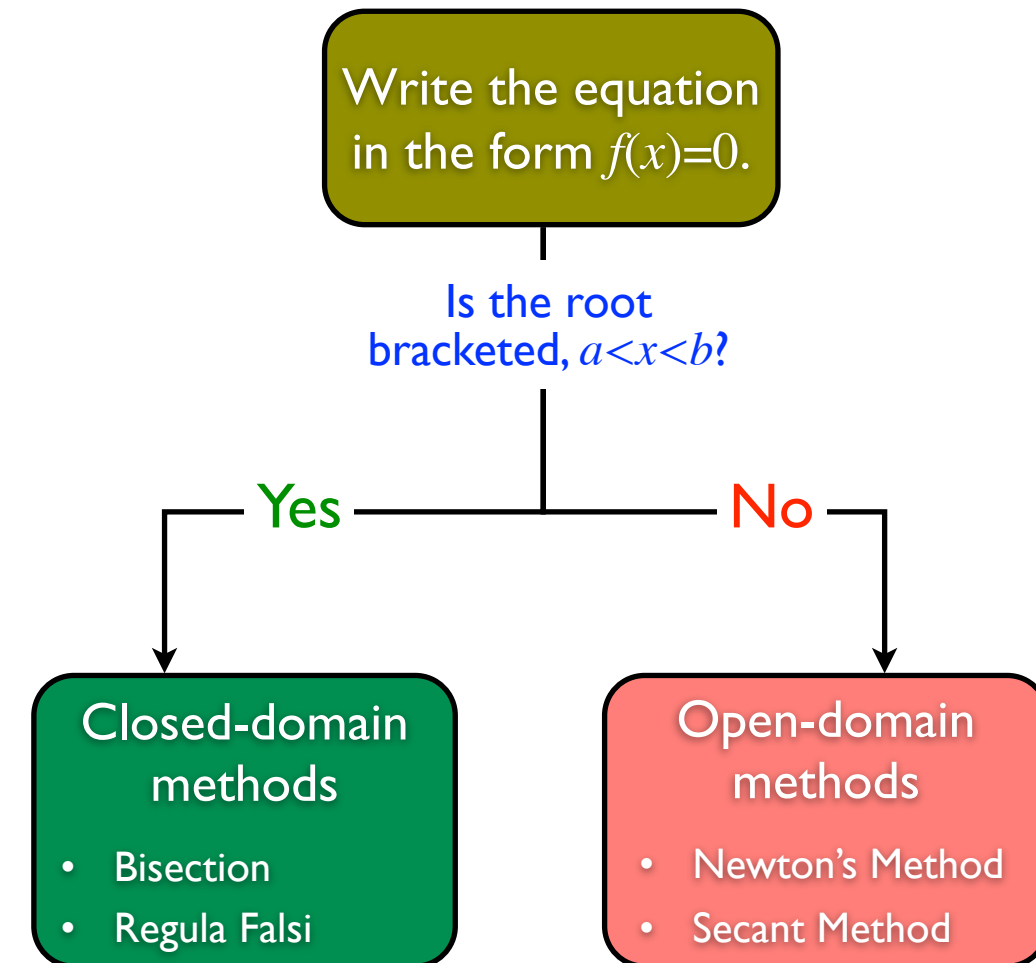
Characteristics:

- May have 0 ... many solutions
- Solution methods are *iterative* and require an *initial guess* for the solution.
- Not guaranteed to find the solution, even if one exists!
 - ▶ Initial guess can be critical to finding the solution.
 - ▶ Bad initial guess may lead to no convergence, or convergence to a wrong (unintended) root.
- Solve for *roots*, $f(x)=0$. If you want $f(x)=a$, then write in residual form, $r(x)=f(x)-a$ and solve $r(x)=0$.



Solution Approaches

- Closed-Domain Methods
 - ▶ Bracket the root and “home in” on it.
 - ▶ Quite simple & robust, but require you to bound the root.
 - ▶ Can be problematic if you bound multiple roots...
- Open Domain Methods
 - ▶ Require an initial guess for the solution, but not a bracket.
 - ▶ More efficient, but less robust than closed-domain methods.



Nonlinear Systems - Newton's Method

$$\mathbf{f}(\mathbf{x}) = 0 \quad \longleftrightarrow \quad \begin{pmatrix} f_1(x_1, x_2, \dots, x_n) \\ f_2(x_1, x_2, \dots, x_n) \\ \vdots \\ f_n(x_1, x_2, \dots, x_n) \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ \vdots \\ 0 \end{pmatrix}$$

\mathbf{x} is a vector of unknowns.

Taylor Series expansion of f_i in terms of \mathbf{x} :

$$f_i(\mathbf{x}) \approx f_i(\mathbf{x}_0) + \sum_{j=1}^n \underbrace{\frac{\partial f_i}{\partial x_j}}_{J_{ij}} \underbrace{(x_j - x_{j0})}_{\Delta x_j} + \mathcal{O}(\Delta x^2)$$

Example: $n=2$ equations:

$$\begin{aligned} f_1(x_1, x_2) &= f_1(x_{1,0}, x_{2,0}) + \frac{\partial f_1}{\partial x_1}(x_1 - x_{1,0}) + \frac{\partial f_1}{\partial x_2}(x_2 - x_{2,0}) \\ f_2(x_1, x_2) &= f_2(x_{1,0}, x_{2,0}) + \frac{\partial f_2}{\partial x_1}(x_1 - x_{1,0}) + \frac{\partial f_2}{\partial x_2}(x_2 - x_{2,0}) \end{aligned}$$

$$\mathbf{f}(\mathbf{x}) \approx \mathbf{f}(\mathbf{x}_0) + \mathbf{J}\Delta\mathbf{x} \quad \Delta\mathbf{x} \text{ is a vector of corrections (updates).}$$

$$\mathbf{J} = \begin{bmatrix} \frac{\partial f_1(\mathbf{x}_0)}{\partial x_1} & \frac{\partial f_1(\mathbf{x}_0)}{\partial x_2} & \dots & \frac{\partial f_1(\mathbf{x}_0)}{\partial x_n} \\ \frac{\partial f_2(\mathbf{x}_0)}{\partial x_1} & \frac{\partial f_2(\mathbf{x}_0)}{\partial x_2} & \dots & \frac{\partial f_2(\mathbf{x}_0)}{\partial x_n} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{\partial f_n(\mathbf{x}_0)}{\partial x_1} & \frac{\partial f_n(\mathbf{x}_0)}{\partial x_2} & \dots & \frac{\partial f_n(\mathbf{x}_0)}{\partial x_n} \end{bmatrix}$$

We want solution at $f(\mathbf{x})=0$. Therefore:

Newton's
Method

$$\mathbf{J}\Delta\mathbf{x} = -\mathbf{f}(\mathbf{x}_0)$$

This is a
linear system
of equations!

Algorithm

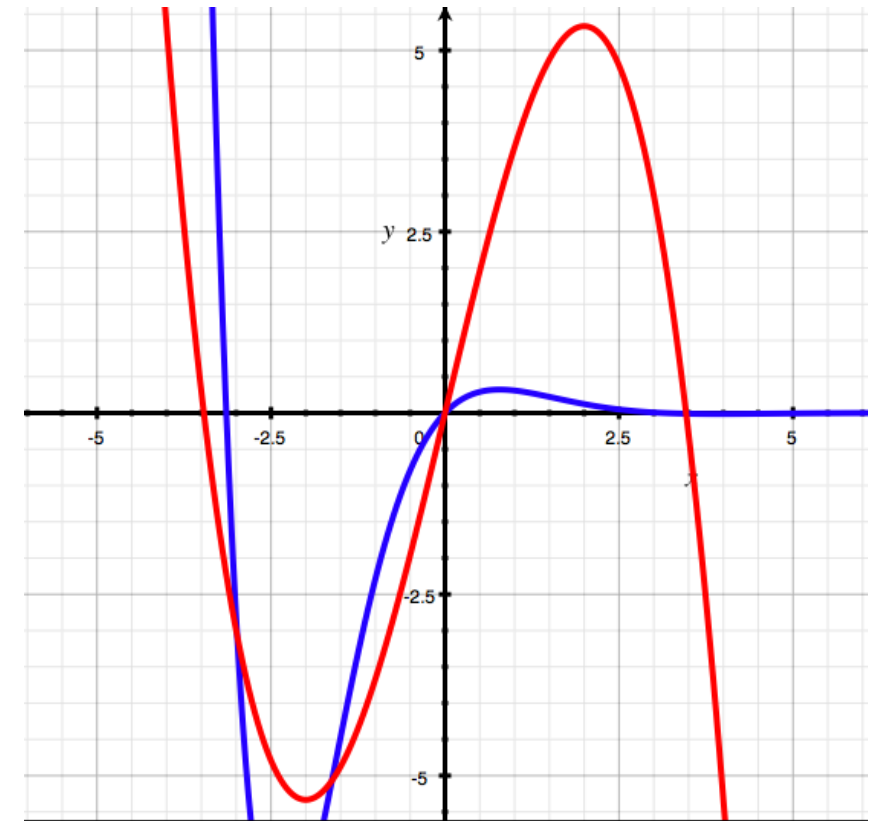
Given $f(\mathbf{x})$, \mathbf{x}_0 , \mathbf{J} .

1. Calculate $[\mathbf{J}]$ & $f(\mathbf{x})$ at the current guess for (\mathbf{x}) .
2. Solve for $(\Delta\mathbf{x})$
3. Update x_i
4. If not converged, go to 1.

Newton's Method - Example

Original Equations: $\frac{1}{2}x^3 + y = 4x$
 $y = \sin(x) \exp(-x)$

Modified Equations: $f_1 = \frac{1}{2}x^3 + y - 4x$
 $f_2 = y - \sin(x) \exp(-x)$



Jacobian: $[J] = \begin{bmatrix} \frac{\partial f_1}{\partial x} & \frac{\partial f_1}{\partial y} \\ \frac{\partial f_2}{\partial x} & \frac{\partial f_2}{\partial y} \end{bmatrix} = \begin{bmatrix} \frac{3}{2}x^2 - 4 & 1 \\ -\cos(x) \exp(-x) + \sin(x) \exp(-x) & 1 \end{bmatrix}$

Example - cont'd

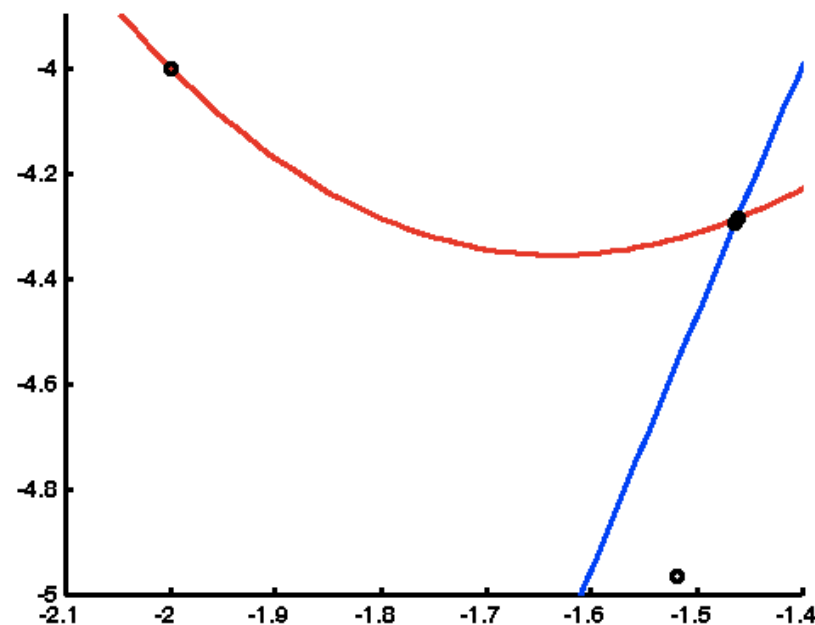
1. Guess x_i : $x = -2, \quad y = -4$

2. Calculate $[J]$ & (f) $[J] = \begin{bmatrix} 2.0 & 1 \\ -3.6439 & 1 \end{bmatrix} \quad (f) = \begin{pmatrix} 0.0 \\ 2.7188 \end{pmatrix}$

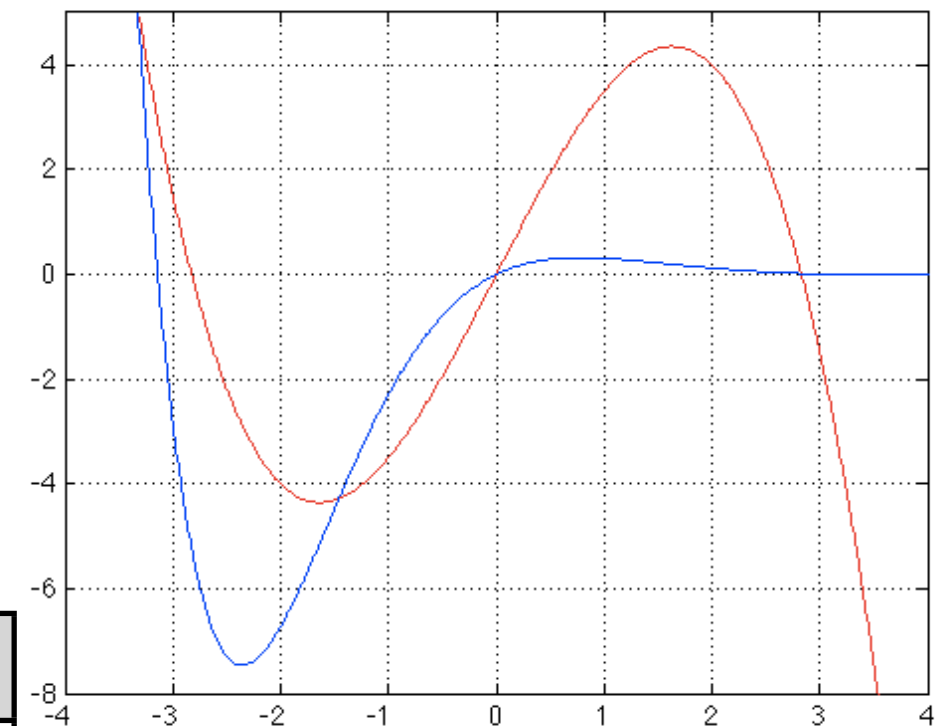
3. Solve for (Δ) $(\Delta) = \begin{pmatrix} 0.4817 \\ -0.9635 \end{pmatrix}$

4. Update x_i $x = -1.5183, \quad y = -4.9635$

$$\begin{aligned} f_1 &= \frac{1}{2}x^3 + y - 4x \\ f_2 &= y - \sin(x) \exp(-x) \end{aligned}$$



k	x	y
0	-2	-4
1	-1.5183	-4.9635
2	-1.4631	-4.2932
3	-1.4611	-4.2848
4	-1.4611	-4.2848



Software Tools

Must provide a function to evaluate the residual.

MATLAB

- FZERO - good for single nonlinear equation, solves for x such that $f(x)=0$.
- FSOLVE - for systems of nonlinear equations, finds x_i such that $f_j(x_i)=0$.
 - requires the “optimization toolbox”
- FMINSEARCH - good for systems of nonlinear equations
 - Searches for the minimum, not the zeros.

Excel

- Goal Seek - single variable
- Solver - multiple variables

Solve the last problem again in MATLAB...

Analysis of PDEs

3-D rectangular coordinate system: $\nabla = \vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y} + \vec{k} \frac{\partial}{\partial z}$

PDEs

$$\underbrace{\rho c_p \frac{\partial T}{\partial t}}_{\text{Time-rate of change in } T \text{ at a point in space.}} = \underbrace{-\rho c_p \mathbf{u} \cdot \nabla T}_{\text{Convection of } T \text{ (velocity pushing it around)}} + \underbrace{\frac{\partial p}{\partial t}}_{T \text{ changes due to changes in } p.} + \underbrace{\mathbf{u} \cdot \nabla p}_{T \text{ changes due to viscous heating.}} + \underbrace{\boldsymbol{\tau} : \nabla \mathbf{u}}_{T \text{ changes due to thermal \& species diffusion.}} - \underbrace{\nabla \cdot \mathbf{q}}_{T \text{ changes from other sources (reaction, radiation, etc).}} + s_T$$

Solution

Assume:

1. Velocity is zero.
2. Pressure is constant.
3. Steady-state.
4. $\mathbf{q} = -\lambda \nabla T$ (Fourier's Law of conduction)
5. λ is constant.
6. One-dimensional

Assume:

1. Velocity is zero.
2. Pressure is constant.
3. T does not vary spatially.
4. $s_T = -hA/V (T - T_\infty)$

$$\frac{d^2 T}{dx^2} = -\frac{s_T}{\lambda}$$

$$\frac{dT}{dt} = -\frac{hA}{\rho c_p V} (T - T_\infty)$$

ODEs

Assumptions

Solution

Algebraic Equations

Numerical Solutions to PDEs

 For systems which have a time derivative ($\partial \phi / \partial t$)

- Convert the PDE into a system of ODEs

- ▶ Method of Lines: “Discretize” in space. Then we are left with a system of ODEs.
- ▶ Number of ODEs is dependent on spatial discretization.

$$\frac{\partial \phi}{\partial t} = \Gamma \frac{\partial^2 \phi}{\partial x^2} + s_\phi$$

“Parabolic” PDE

 For PDEs which do not have a time derivative (Elliptic PDEs):

- Called “boundary value problems”
- Convert to a big system of (nonlinear) equations.
- Number of equations depends on spatial discretization (next).

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = -\frac{s_\phi}{D_\phi}$$

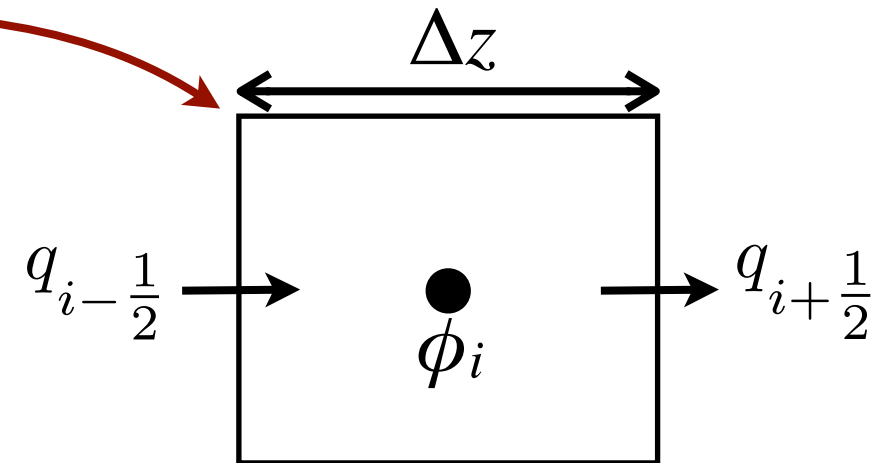
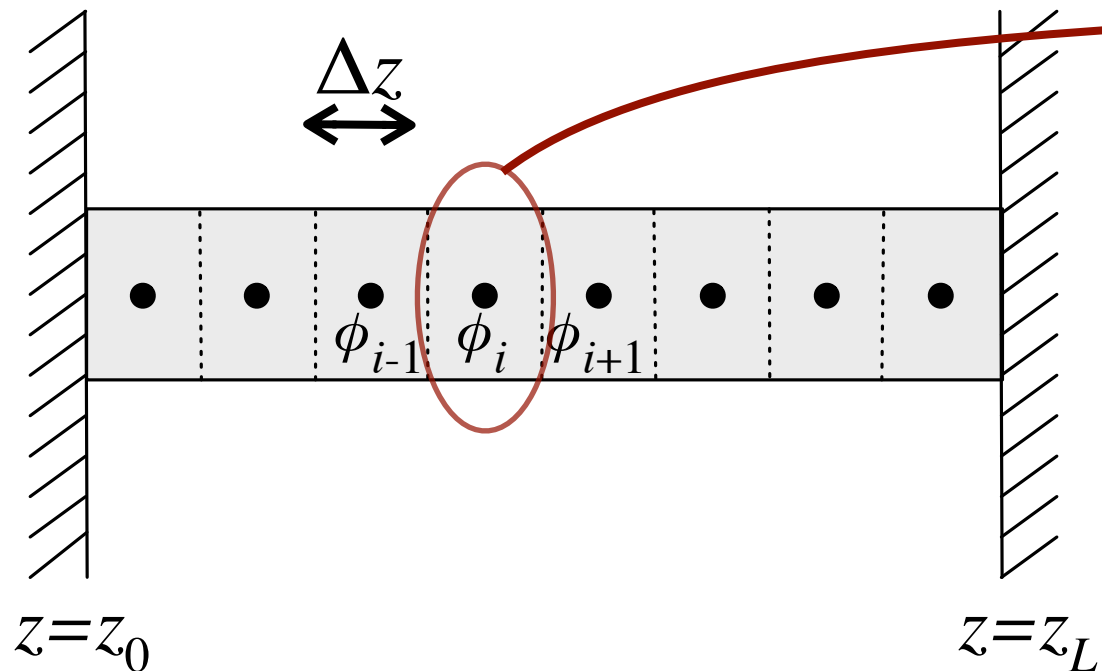
Elliptic PDEs

Here we will show examples primarily for
Boundary-Value ODEs

Elliptic PDE: $\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = -\frac{s_\phi}{D_\phi}$

Boundary
value ODE: $\frac{d^2 \phi}{dx^2} = -\frac{s_\phi}{D_\phi}$

Discrete Calculus (I-D)



Assume a flux, q , of the form:

$$q = -D \frac{d\phi}{dz}$$

We may approximate q as:

$$q_{i+\frac{1}{2}} \approx -D_{i+\frac{1}{2}} \frac{\phi_{i+1} - \phi_i}{\Delta z}$$

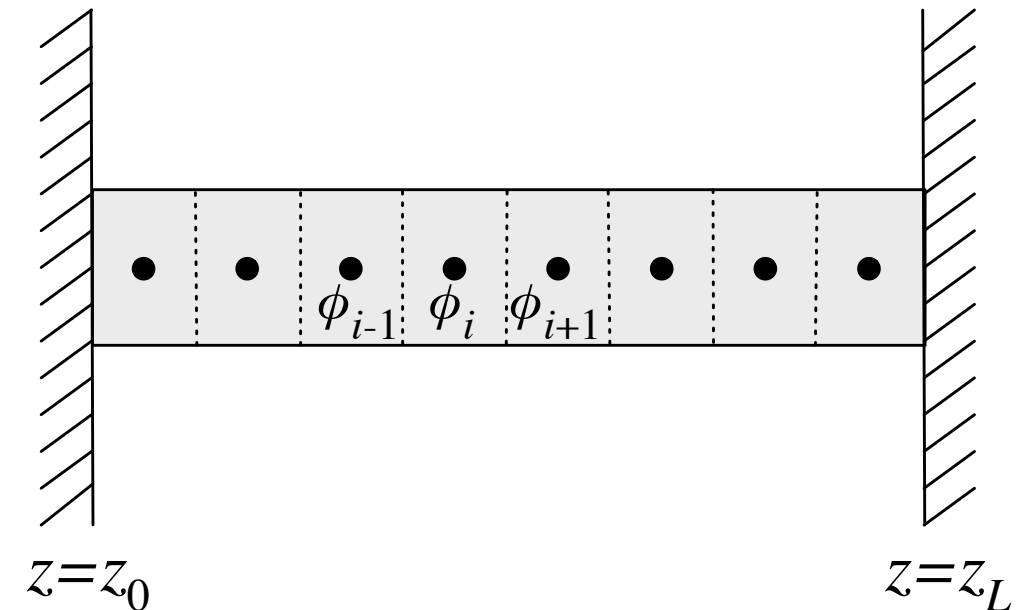
$$q_{i-\frac{1}{2}} \approx -D_{i-\frac{1}{2}} \frac{\phi_i - \phi_{i-1}}{\Delta z}$$

$$\left. \frac{d\phi}{dz} \right|_{i+\frac{1}{2}} = \frac{\phi_{i+1} - \phi_i}{\Delta z} + \mathcal{O}(\Delta z^2)$$

Approximation for the derivative of ϕ at the midpoint of two points.

Second Derivatives

$$\left. \frac{d\phi}{dz} \right|_{i+\frac{1}{2}} = \frac{\phi_{i+1} - \phi_i}{\Delta z} + \mathcal{O}(\Delta z^2)$$



Use the same approximation on the *derivatives* of ϕ to obtain:

$$\begin{aligned} \left. \frac{d^2\phi}{dz^2} \right|_i &= \frac{\left. \frac{d\phi}{dz} \right|_{i+\frac{1}{2}} - \left. \frac{d\phi}{dz} \right|_{i-\frac{1}{2}}}{\Delta z} + \mathcal{O}(\Delta z^2), \\ &= \frac{1}{\Delta z} \left[\frac{\phi_{i+1} - \phi_i}{\Delta z} - \frac{\phi_i - \phi_{i-1}}{\Delta z} \right] + \mathcal{O}(\Delta z^2) \end{aligned}$$

$$\left. \frac{d^2\phi}{dz^2} \right|_i = \frac{\phi_{i+1} - 2\phi_i + \phi_{i-1}}{\Delta z^2} + \mathcal{O}(\Delta z^2)$$

Approximation for the second derivative, valid for uniformly spaced cells.

Example - Steady Diffusion

Steady state, no convection: $\nabla \cdot \mathbf{q} = s$

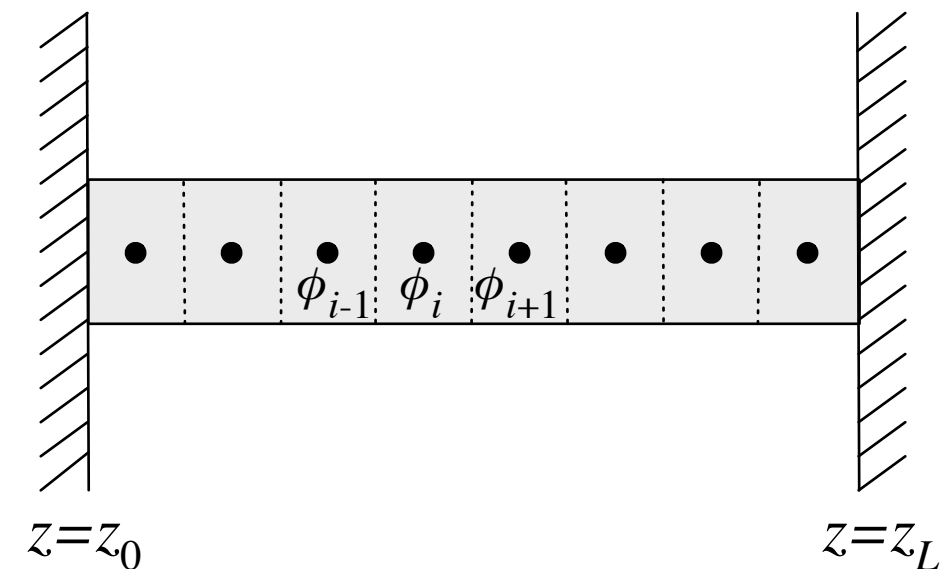
$$\nabla \cdot \mathbf{J}_i = \frac{s_i}{M_i}$$

“Effective binary” or heat conduction: $\mathbf{q} = -D\nabla\phi$

Constant diffusivity: $\nabla^2\phi = -\frac{s}{D}$

One-dimensional: $\frac{d^2\phi}{dz^2} = -\frac{s}{D}$

$$\left. \frac{d^2\phi}{dz^2} \right|_i = \frac{\phi_{i+1} - 2\phi_i + \phi_{i-1}}{\Delta z^2} + \mathcal{O}(\Delta z^2)$$



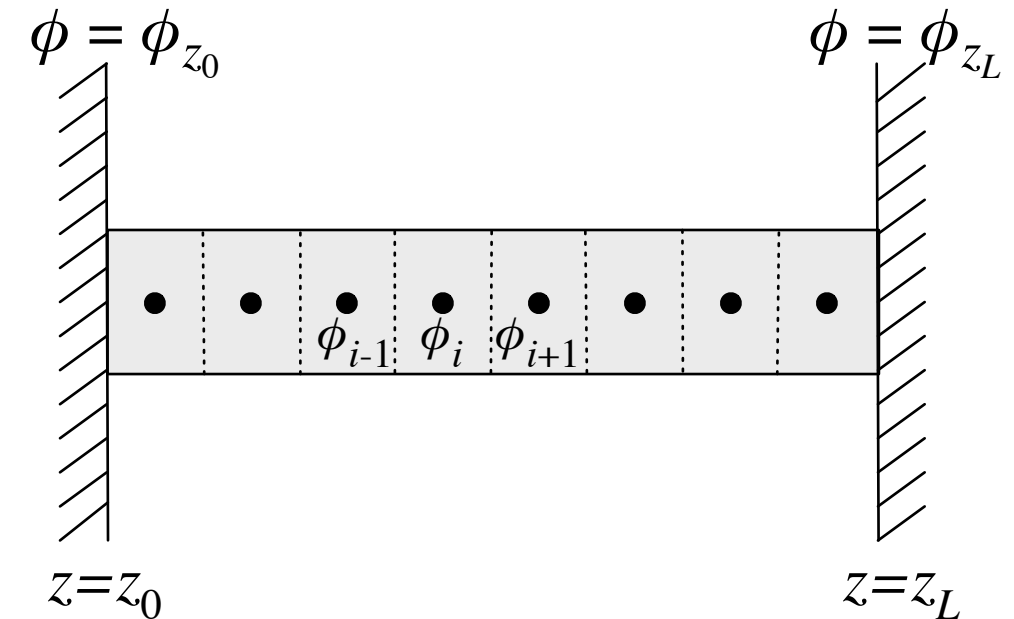
$$\frac{\phi_{i+1} - 2\phi_i + \phi_{i-1}}{\Delta z^2} = -\frac{s_i}{D}$$

We can apply this at all “interior” points. At the boundaries, we must modify this....

Dirichlet Boundary Conditions

If the solution variable is known at the boundary, then we call this a **Dirichlet** boundary condition.

$$\frac{\phi_{i+1} - 2\phi_i + \phi_{i-1}}{\Delta z^2} = -\frac{s_i}{D} \quad \text{This applies at all interior points, } 2 \leq i \leq n-1.$$



Linear interpolation
between "0" and "1"

At z_0 , $\phi = \phi_{z_0}$. $\frac{\phi_0 + \phi_1}{2} = \phi_{z_0} \Rightarrow \phi_0 = 2\phi_{z_0} - \phi_1$

Using the top equation, $\frac{\phi_2 - 3\phi_1}{\Delta z^2} = -\frac{s_1}{D} - \frac{2\phi_{z_0}}{\Delta z^2}$ applies at $i=1$.

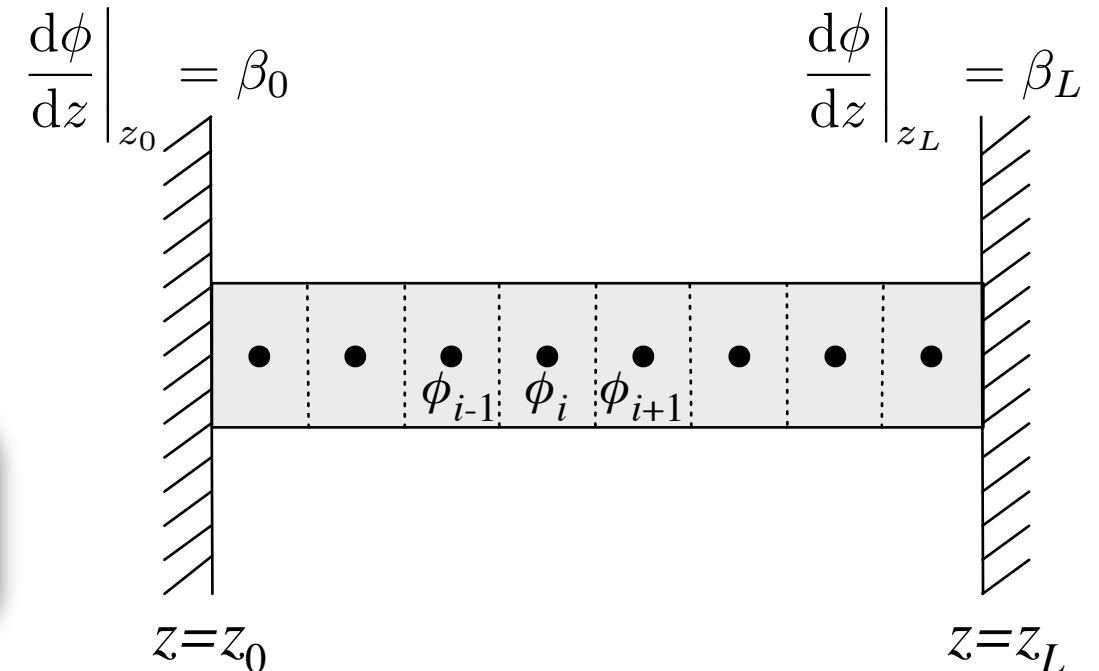
At z_L , $\phi = \phi_{z_L}$. $\frac{\phi_{n+1} + \phi_n}{2} = \phi_{z_L} \Rightarrow \phi_{n+1} = 2\phi_{z_L} - \phi_n$

Using the top equation, $\frac{\phi_{n-1} - 3\phi_n}{\Delta z^2} = -\frac{s_n}{D} - \frac{2\phi_{z_L}}{\Delta z^2}$ applies at $i=n$.

Neumann Boundary Conditions

If the derivative solution variable is known at the boundary, then we call this a **Neumann** boundary condition.

$$\frac{\phi_{i+1} - 2\phi_i + \phi_{i-1}}{\Delta z^2} = -\frac{s_i}{D} \quad \text{This applies at all interior points, } 2 \leq i \leq n-1.$$



At $i=1$, $\frac{\phi_1 - \phi_0}{\Delta z} = \beta_0 \Rightarrow \phi_0 = \phi_1 - \beta_0 \Delta z$

Using the top equation,

$$\frac{-\phi_1 + \phi_2}{\Delta z^2} = -\frac{s_1}{D} - \frac{\beta_0}{\Delta z} \quad \text{applies at } i=1.$$

At $i=n$, $\frac{\phi_{n+1} - \phi_n}{\Delta z} = \beta_L \Rightarrow \phi_{n+1} = \beta_L \Delta z + \phi_n$

Using the top equation,

$$\frac{-\phi_n + \phi_{n-1}}{\Delta z^2} = -\frac{s_n}{D} - \frac{\beta_L}{\Delta z} \quad \text{applies at } i=n.$$

Example: Steady Conduction

If it were species rather than temperature, then this looks like a “Diffusion-reaction balance”

$$\frac{d^2 T}{dz^2} = -\frac{s(z)}{\lambda} \quad s = \exp\left(-\frac{\left(z - \frac{L}{2}\right)^2}{\gamma}\right)$$

Boundary conditions:

$$T(z = 0) = 0$$

$$\left. \frac{dT}{dz} \right|_{z=L} = 0$$

$$\lambda = 10^{-3}$$

$$\gamma = 10^{-3}$$

$$L = 1$$

Find $T(z)$.

Steps:

1. Write down the discrete equations for interior
2. Write discrete equations at the boundary.
3. Write the matrix to be solved.
4. Finally, go to Matlab to solve the problem.

Interior equations $(1 < i < n)$
$$\frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta z^2} = -\frac{s_i}{\lambda}$$

$$s_i = \exp\left(-\frac{\left(z_i - \frac{L}{2}\right)^2}{\gamma}\right)$$

Left Boundary $(i = 1)$
$$\frac{T_0 - 2T_1 + T_2}{\Delta z^2} = -\frac{s_1}{\lambda}$$
 (must eliminate T_0)
$$\frac{T_1 + T_0}{2} = T(z = 0) = 0$$

$$\frac{-3T_1 + T_2}{\Delta z^2} = -\frac{s_1}{\lambda}$$

Right Boundary $(i = n)$
$$\frac{T_{n-1} - 2T_n + T_{n+1}}{\Delta z^2} = -\frac{s_n}{\lambda}$$
 (must eliminate T_{n+1})
$$\frac{T_{n+1} - T_n}{\Delta z} = \frac{dT}{dz}\bigg|_{z=L} = 0$$

$$\frac{T_{n-1} - T_n}{\Delta z^2} = -\frac{s_n}{\lambda}$$

Left boundary condition

For 5 control volumes, we have:

$$\begin{bmatrix} -3 & 1 & 0 & 0 & 0 \\ 1 & -2 & 1 & 0 & 0 \\ 0 & 1 & -2 & 1 & 0 \\ 0 & 0 & 1 & -2 & 1 \\ 0 & 0 & 0 & 1 & -1 \end{bmatrix} \begin{pmatrix} T_1 \\ T_2 \\ T_3 \\ T_4 \\ T_5 \end{pmatrix} = -\frac{\Delta z^2}{\lambda} \begin{pmatrix} s_1 \\ s_2 \\ s_3 \\ s_4 \\ s_5 \end{pmatrix}$$

Right boundary condition

Nonlinear BVPs

Example: $\frac{d^2 T}{dx^2} = -\alpha (T^4 - T_\infty^4)$

$$\left. \frac{d^2 T}{dx^2} \right|_i \approx \frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta x^2}$$

Nonlinear
term

Discrete equation to solve
at each “interior” point $\frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta x^2} = -\alpha (T_i^4 - T_\infty^4)$

Options:

1. Leave T_i^4 on the right hand side & try to solve the linear system (**not** a good option).
2. Solve the nonlinear system of equations using Newton's method.
 - rewrite in residual form
 - requires a Jacobian matrix
 - This is the most general approach (big hammer)
3. *Linearize* the equation.



Linearization

Example: $y = 5x^3 - 2x$

Taylor series expansion about x_o :

$$f(x) = f(x_o) + f'(x_o)(x - x_o) + \frac{1}{2}f''(x_o)(x - x_o)^2 + \cdots + \frac{1}{n!}f^{(n)}(x_o)(x - x_o)^{n+1}$$

$$\begin{aligned} y &\approx 5x_o^3 - 2x_o + (15x_o^2 - 2)(x - x_o) \\ &= -10x_o^3 + (15x_o^2 - 2)x \end{aligned}$$

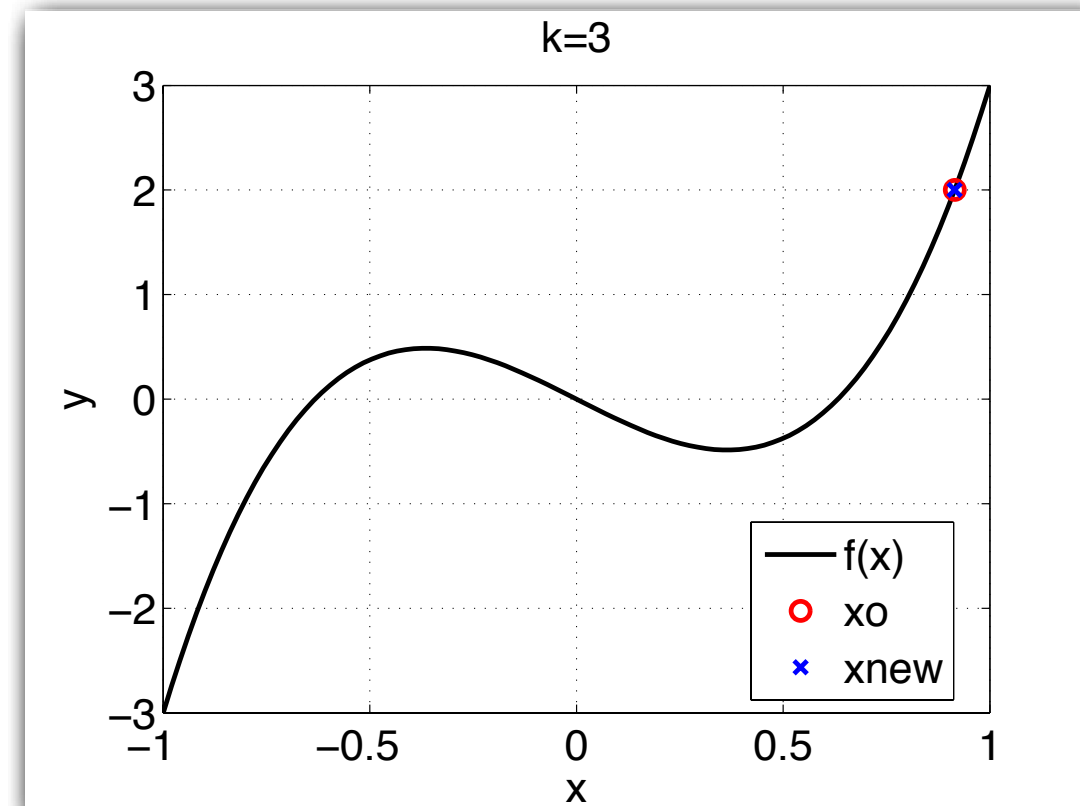
Now y is *linear* with respect to x
(nonlinear with respect to x_o).

Example: Solve for x such that $y=2$. $x = \frac{y + 10x_o^3}{15x_o^2 - 2}$

1. Guess x_o .
2. Calculate new value for x .
3. if $|x - x_o| > \varepsilon$ then $x_o = x$, return to step 2. Otherwise, done!

k	x_o	x
1	1	0.9231
2	0.9231	0.9151
3	0.9151	0.9150

Exercise: what happens when we change our initial guess to $x=0$?



Linearization for Nonlinear BVPs

$$\frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta x^2} = -\alpha (T_i^4 - T_\infty^4)$$

Linearize T_i^4 term about T_i^* :

$$T_i^4 \approx (T_i^*)^4 + 4(T_i^*)^3(T_i - T_i^*)$$

$$\frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta x^2} = -\alpha [(T_i^*)^4 + 4(T_i^*)^3(T_i - T_i^*) - T_\infty^4]$$

$$\frac{1}{\Delta x^2} T_{i-1} - \left(\frac{2}{\Delta x^2} + 4\alpha(T_i^*)^3 \right) T_i + \frac{1}{\Delta x^2} T_{i+1} = \alpha [3(T_i^*)^4 + T_\infty^4]$$

Applies to all interior points.

$$\begin{bmatrix} BC_1 & & & \\ \frac{1}{\Delta x^2} & -\left(\frac{2}{\Delta x^2} + 4\alpha(T_2^*)^3\right) & \frac{1}{\Delta x^2} & 0 \\ \ddots & \ddots & \ddots & \\ 0 & \frac{1}{\Delta x^2} & -\left(\frac{2}{\Delta x^2} + 4\alpha(T_{n-1}^*)^3\right) & \frac{1}{\Delta x^2} \\ & & & BC_n \end{bmatrix} \begin{pmatrix} T_1 \\ T_2 \\ \vdots \\ T_{n-1} \\ T_n \end{pmatrix} = \begin{pmatrix} bc_1 \\ \alpha [3(T_2^*)^4 + T_\infty^4] \\ \vdots \\ \alpha [3(T_{n-1}^*)^4 + T_\infty^4] \\ bc_n \end{pmatrix}$$

Boundary conditions implemented as previously discussed.

1. Guess the solution values (T_i^*)
2. Update the LHS matrix and RHS vector given these values for T_i^* .
3. Solve the system of equations for T_i .
4. If $\|T_i - T_i^*\| > \varepsilon$ then set $T_i^* = T_i$ and return to step 2. Otherwise, we have the answer.

You choose the norm you want (L_2, L_∞)

“Elliptic” PDEs

In chemical engineering applications, elliptic PDEs typically arise from steady-state diffusion problems.

$$\nabla^2 \phi = f(\vec{x}, \phi)$$

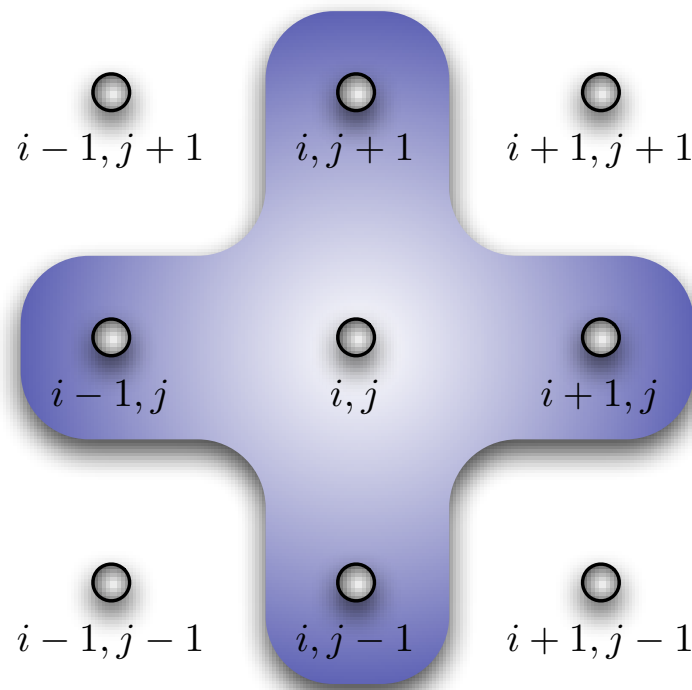
$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = f(x, y, \phi)$$

two-dimensional
rectangular coordinates

Second-order discretization:

$$\frac{\phi_{i-1,j} - 2\phi_{i,j} + \phi_{i+1,j}}{\Delta x^2} + \frac{\phi_{i,j-1} - 2\phi_{i,j} + \phi_{i,j+1}}{\Delta y^2} = f(x_{i,j}, y_{i,j}, \phi_{i,j})$$

Applies at all “interior” points.



- At $i=1$, and $i=n_x$ apply x boundary conditions.
- At $j=1$, and $j=n_y$ apply y boundary conditions.

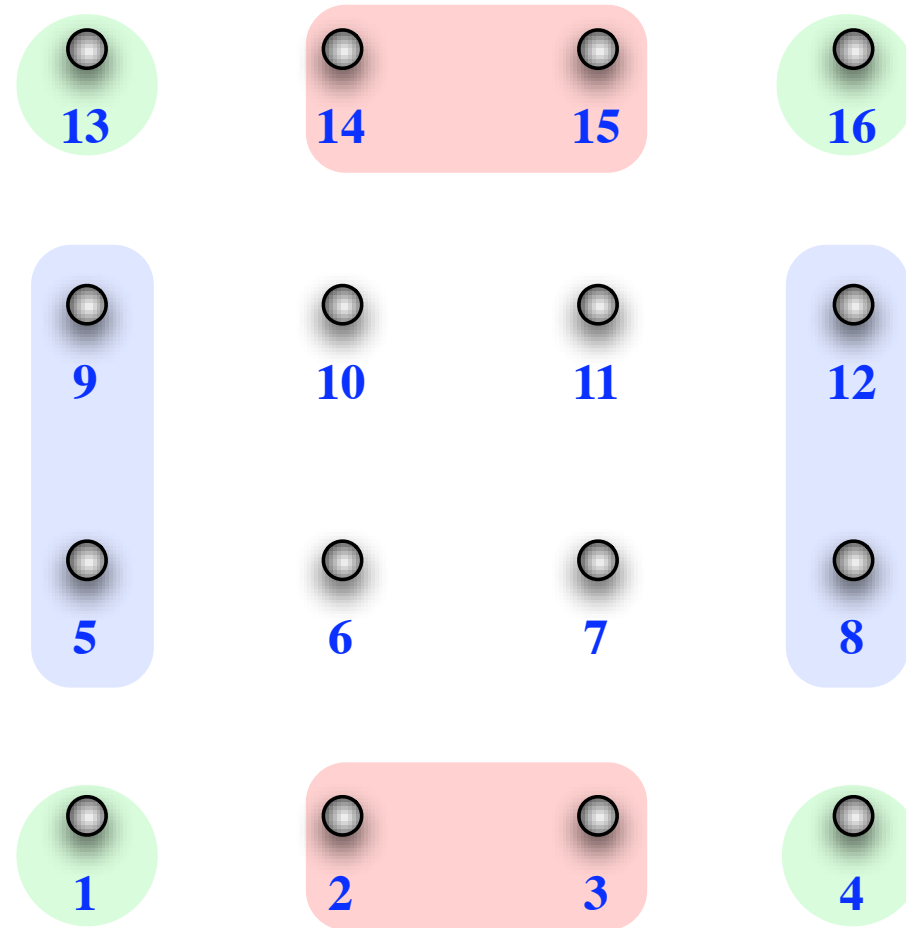
Variable numbering (4x4 grid)

- (x,y) layout
- solution index (eqn #) layout

1,4	2,4	3,4	4,4
13	14	15	16
1,3	2,3	3,3	4,3
9	10	11	12
1,2	2,2	3,2	4,2
5	6	7	8
1,1	2,1	3,1	4,1
1	2	3	4

$$\frac{\phi_{i-1,j} - 2\phi_{i,j} + \phi_{i+1,j}}{\Delta x^2} + \frac{\phi_{i,j-1} - 2\phi_{i,j} + \phi_{i,j+1}}{\Delta y^2} = f(x_{i,j}, y_{i,j}, \phi_{i,j})$$

Note: if $f(x,y)$ depends on ϕ then this is a system of **nonlinear** equations!



BC	??	0	0	??	0	0	0	0	0	0	0	0	0	0	0	0	0	$\begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_4 \\ \phi_5 \\ \phi_6 \\ \phi_7 \\ \phi_8 \\ \phi_9 \\ \phi_{10} \\ \phi_{11} \\ \phi_{12} \\ \phi_{13} \\ \phi_{14} \\ \phi_{15} \\ \phi_{16} \end{pmatrix} = \begin{pmatrix} bc_1 \\ f(x_2, y_2, \phi_2) + bc_2 \\ f(x_3, y_3, \phi_3) + bc_3 \\ bc_4 \\ f(x_5, y_5, \phi_5) + bc_5 \\ f(x_6, y_6, \phi_6) \\ f(x_7, y_7, \phi_7) \\ f(x_8, y_8, \phi_8) + bc_8 \\ f(x_9, y_9, \phi_9) + bc_9 \\ f(x_{10}, y_{10}, \phi_{10}) \\ f(x_{11}, y_{11}, \phi_{11}) \\ f(x_{12}, y_{12}, \phi_{12}) + bc_{12} \\ f(x_7, y_7, \phi_7) + bc_{13} \\ f(x_{14}, y_{14}, \phi_{14}) + bc_{14} \\ f(x_{15}, y_{15}, \phi_{15}) + bc_{15} \\ bc_{16} \end{pmatrix}$
$\frac{1}{\Delta x^2}$	$-\frac{2}{\Delta x^2}$	$\frac{1}{\Delta x^2}$	0	0	??	0	0	0	0	0	0	0	0	0	0	0	0	
0	$\frac{1}{\Delta x^2}$	$-\frac{1}{\Delta x^2}$	$\frac{1}{\Delta x^2}$	0	0	??	0	0	0	0	0	0	0	0	0	0	0	
0	0	??	BC	0	0	??	0	0	0	0	0	0	0	0	0	0	0	
$\frac{1}{\Delta y^2}$	0	0	0	$-\frac{2}{\Delta x^2}$??	0	0	$\frac{1}{\Delta y^2}$	0	0	0	0	0	0	0	0	0	
0	$\frac{1}{\Delta y^2}$	0	0	$\frac{1}{\Delta x^2}$	$-2\left(\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2}\right)$	$\frac{1}{\Delta x^2}$	0	0	$\frac{1}{\Delta y^2}$	0	0	0	0	0	0	0	0	
0	0	$\frac{1}{\Delta y^2}$	0	0	$\frac{1}{\Delta x^2}$	$-2\left(\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2}\right)$	$\frac{1}{\Delta x^2}$	0	0	$\frac{1}{\Delta y^2}$	0	0	0	0	0	0	0	
0	0	0	$\frac{1}{\Delta y^2}$	0	0	??	$-\frac{2}{\Delta y^2}$	0	0	0	$\frac{1}{\Delta y^2}$	0	0	0	0	0	0	
0	0	0	0	$\frac{1}{\Delta y^2}$	0	0	0	$-\frac{2}{\Delta y^2}$??	0	0	$\frac{1}{\Delta y^2}$	0	0	0	0	0	
0	0	0	0	0	$\frac{1}{\Delta y^2}$	0	0	$\frac{1}{\Delta x^2}$	$-2\left(\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2}\right)$	$\frac{1}{\Delta x^2}$	0	0	$\frac{1}{\Delta y^2}$	0	0	0		
0	0	0	0	0	0	$\frac{1}{\Delta y^2}$	0	0	$\frac{1}{\Delta x^2}$	$-2\left(\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2}\right)$	$\frac{1}{\Delta x^2}$	0	0	$\frac{1}{\Delta y^2}$	0	0		
0	0	0	0	0	0	0	$\frac{1}{\Delta y^2}$	0	0	??	$-\frac{2}{\Delta y^2}$	0	0	0	0	$\frac{1}{\Delta y^2}$	0	
0	0	0	0	0	0	0	0	0	??	0	0	BC	??	0	0	0	0	
0	0	0	0	0	0	0	0	0	0	0	0	$\frac{1}{\Delta x^2}$	$-\frac{2}{\Delta x^2}$	$\frac{1}{\Delta x^2}$	0	0		
0	0	0	0	0	0	0	0	0	0	0	0	0	$\frac{1}{\Delta x^2}$	$-\frac{2}{\Delta x^2}$	$\frac{1}{\Delta x^2}$	$\frac{1}{\Delta x^2}$		
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	??	BC	0	

Parabolic PDEs

$$\rho c_p \frac{\partial T}{\partial t} = \lambda \frac{\partial^2 T}{\partial x^2} + s_T$$

Ordinary Differential Equations (ODEs)

A coupled system of ODEs: $\frac{d\phi_i}{dt} = F_i(\phi_j)$

Explicit:

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = F_i(\phi_j^n) + \mathcal{O}(\Delta t)$$

Use for:

- “Non-stiff” equations
- Unsteady solutions

Implicit:

$$\frac{\phi_i^{n+1} - \phi_i^n}{\Delta t} = F_i(\phi_j^{n+1}) + \mathcal{O}(\Delta t)$$



$$r_i = F_i(\phi_j^{n+1}) - \frac{\phi_i^{n+1} - \phi_i^n}{\Delta t}$$

System of nonlinear equations.
Solve for $r_i = 0$.

Use for:

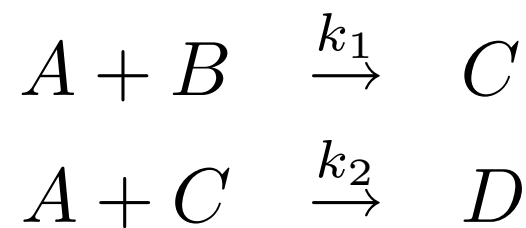
- “Stiff” equations
- Time-marching to steady solutions

Higher-order methods can be constructed. (Runge Kutta, etc.)

MATLAB: **ode45** (nonstiff), **ode23s** (stiff)

You provide evaluation of $F_i(\phi_j)$




Example: Kinetics



$$\begin{array}{l} \frac{dC_A}{dt} = -r_1 - r_2 \\ \frac{dC_B}{dt} = -r_1 \\ \frac{dC_C}{dt} = r_1 - r_2 \\ \frac{dC_D}{dt} = r_2 \end{array}$$

Initial conditions: $C_A=1, C_B=0.6$.

Rate constants: $k_1=1, k_2=0.1$.

-  Plot concentrations as functions of time.
 - requires solution of system of ODEs
-  Determine the equilibrium composition.
 - Use stoichiometry & mole balance
-  How long to achieve 99% of equilibrium?
 - find this entry in the ODE solution history.

“Parabolic” PDEs

Parabolic PDEs are characterized primarily by
transient diffusion.

Transient diffusion equation
(constant diffusivity, Γ)

$$\frac{\partial \phi}{\partial t} = \Gamma \nabla^2 \phi + s_\phi$$

Transient temperature diffusion equation
(constant pressure, diffusivity)

$$\rho c_p \frac{\partial T}{\partial t} = \lambda \nabla^2 T + s_T$$

Transient temperature diffusion equation
(two-dimensional version of the above equation)

$$\rho c_p \frac{\partial T}{\partial t} = \lambda \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) + s_T$$

Solving Time-Dependent PDEs

The “Method of Lines”

$$\frac{\partial \phi}{\partial t} = \Gamma \frac{\partial^2 \phi}{\partial x^2} + s$$

Spatial discretization:

$$\frac{\partial^2 \phi}{\partial x^2} \approx \frac{\phi_{i-1} - 2\phi_i + \phi_{i+1}}{\Delta x^2}$$

$$\frac{\partial \phi_i}{\partial t} = \Gamma_i \frac{\phi_{i-1} - 2\phi_i + \phi_{i+1}}{\Delta x^2} + s_i$$

This is a system of coupled ODEs.

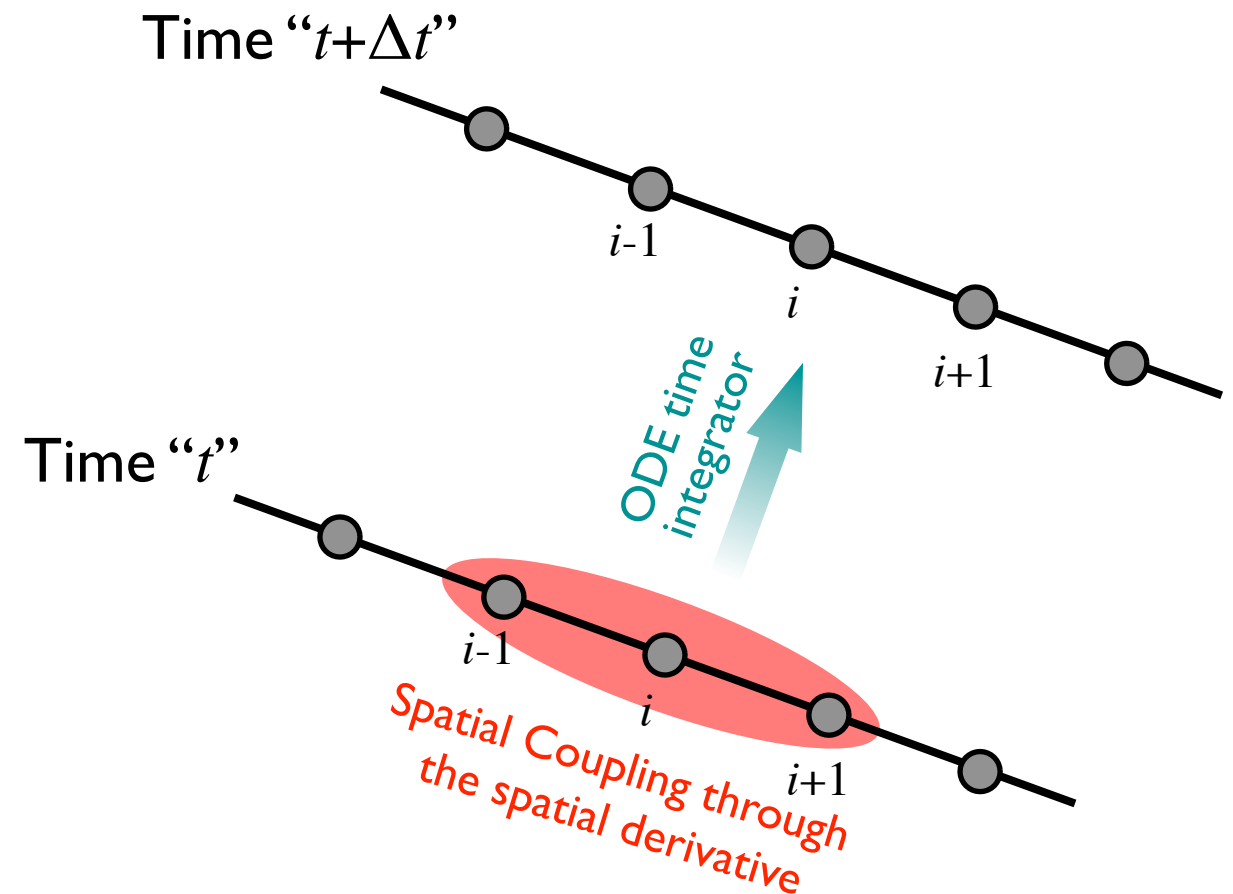
$$\frac{\partial \phi_1}{\partial t} = ? \text{ Need to use BC information...}$$

$$\frac{\partial \phi_2}{\partial t} = \Gamma \frac{\phi_1 - 2\phi_2 + \phi_3}{\Delta x^2} + s_2,$$

$$\vdots = \vdots$$

$$\frac{\partial \phi_{n-1}}{\partial t} = \Gamma \frac{\phi_{n-2} - 2\phi_{n-1} + \phi_n}{\Delta x^2} + s_{n-1},$$

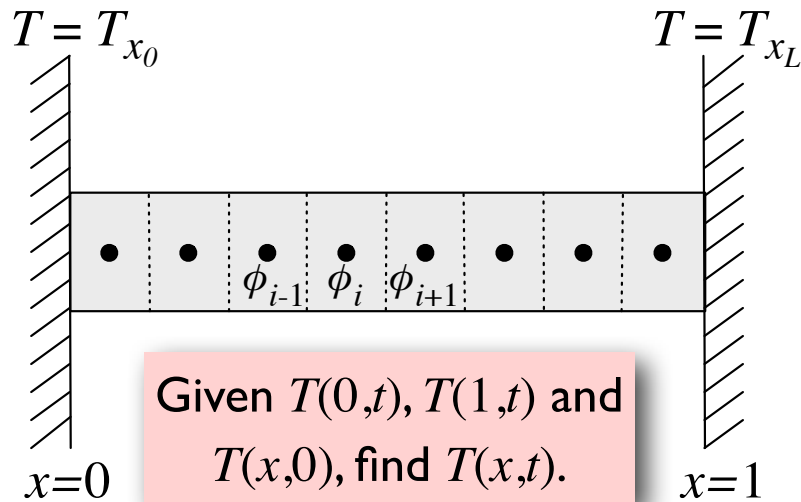
$$\frac{\partial \phi_n}{\partial t} = ? \text{ Need to use BC information...}$$



Solve using Matlab's **ode45** or **ode23s**.

Example I

(Dirichlet Boundary Conditions)



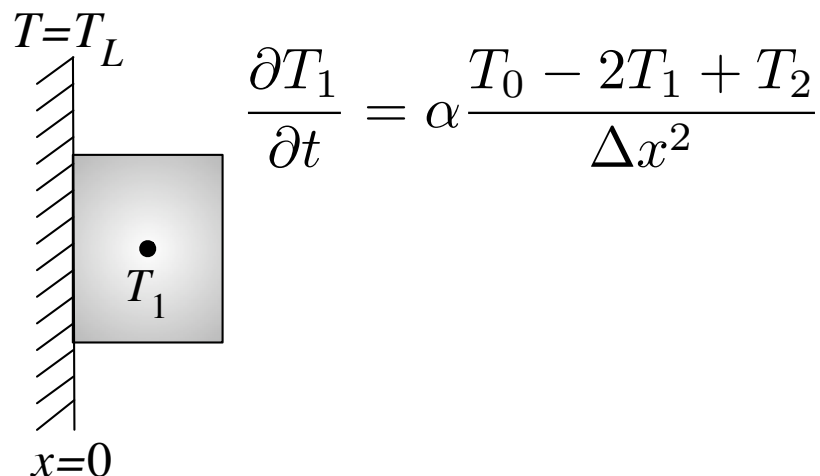
$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

$$\begin{aligned} T(0,t) &= 0, & \text{BC at } x=0 \\ T(1,t) &= 0, & \text{BC at } x=1 \\ T(x,0) &= \phi(x) & \text{Initial condition} \end{aligned}$$

Analytical solution: $T(x,t) = \sum_{k=1}^{\infty} A_k e^{-k^2 \pi^2 \alpha t} \sin(k\pi x)$ $A_k = 2 \int_0^1 \phi(x) \sin(k\pi x) dx$

Numerical solution: $\frac{\partial T_i}{\partial t} = \alpha \frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta x^2}$ Applies at all "interior" points ($i=2, \dots, n-1$)

What do we do at the boundaries ($i=1, i=n$)?



$$T_{x_0} = \frac{T_1 + T_0}{2} \Rightarrow T_0 = 2T_{x_0} - T_1$$

$$\frac{\partial T_1}{\partial t} = \alpha \frac{2T_{x_0} - 3T_1 + T_2}{\Delta x^2}$$

A similar procedure at $x=L$ leads to:

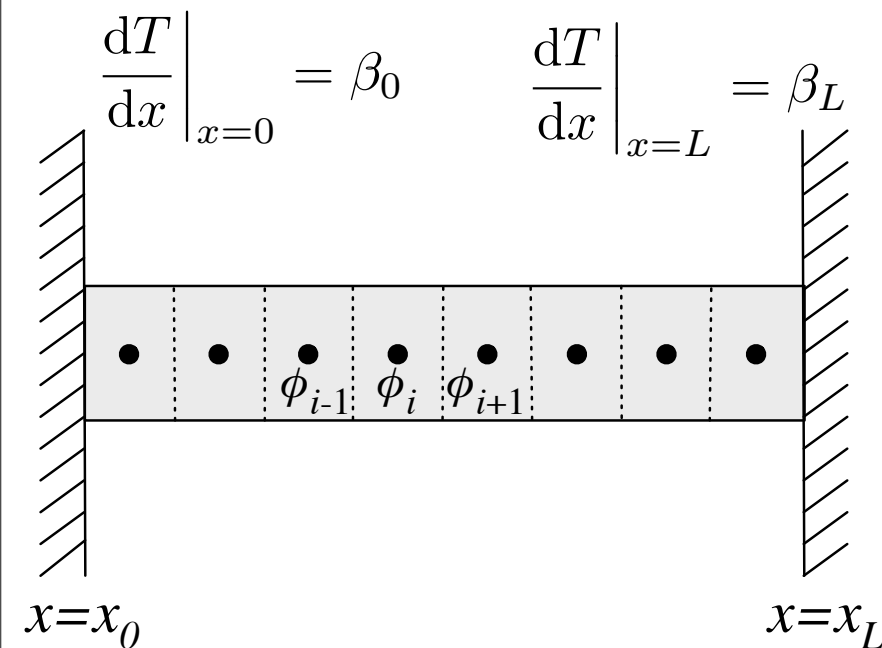
$$\frac{\partial T_n}{\partial t} = \alpha \frac{2T_{x_L} - 3T_n + T_{n-1}}{\Delta x^2}$$

Solve this problem for $t = [0, 240]$ seconds given: $\phi(x) = \exp(-100(x - 0.3)^2)$

$$\alpha = \frac{\lambda}{\rho c_p} = 10^{-4} \frac{\text{m}^2}{\text{s}}$$

Example 2

(Neumann Boundary Conditions)



$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

$$\begin{aligned} \frac{\partial T}{\partial x}\bigg|_{0,t} &= \beta_0, & \text{BC at } x=0 \\ \frac{\partial T}{\partial x}\bigg|_{1,t} &= \beta_L, & \text{BC at } x=1 \\ T(x,0) &= \phi(x) & \text{Initial condition} \end{aligned}$$

Numerical solution: $\frac{\partial T_i}{\partial t} = \alpha \frac{T_{i-1} - 2T_i + T_{i+1}}{\Delta x^2}$ Applies at all "interior" points ($i=2 \dots n-1$)

What do we do at the boundaries ($i=1, i=n$)?

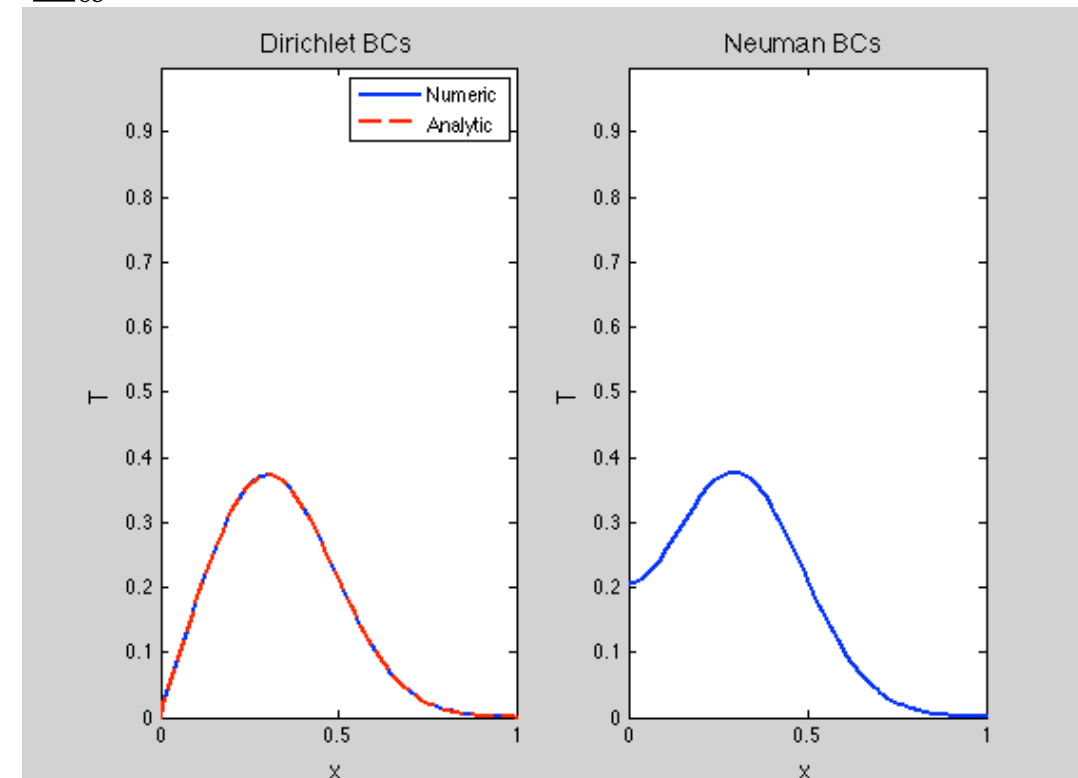
At $x=0$, $\frac{dT}{dx}\big|_{x=0} = \beta_0$. The finite difference approximation is $\frac{\partial T_1}{\partial t} = \alpha \frac{T_0 - 2T_1 + T_2}{\Delta x^2}$. Need to eliminate T_0 . Using $\frac{dT}{dx}\big|_{x=0} = \frac{T_1 - T_0}{\Delta x} = \beta_0 \Rightarrow T_0 = T_1 - \beta_0 \Delta x$.

$$\frac{\partial T_1}{\partial t} = \alpha \left(\frac{T_2 - T_1}{\Delta x^2} - \frac{\beta_0}{\Delta x} \right) \text{ at } i=1.$$

Using the same analysis at $x=1$ ($i=n$)...

$$\frac{\partial T_n}{\partial t} = \alpha \frac{T_{n-1} - T_n}{\Delta x^2} + \frac{\alpha}{\Delta x} \beta_L \text{ at } i=n.$$

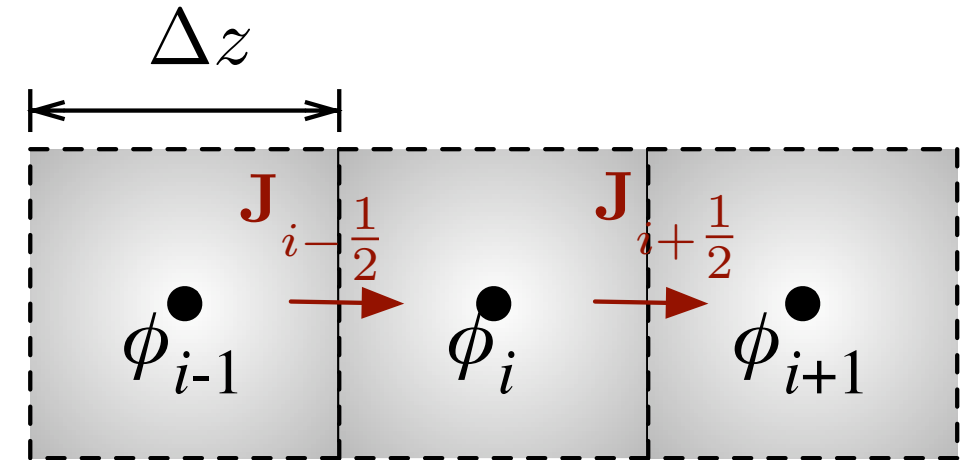
Solve this problem given $\beta_0=\beta_L=0$ and the same conditions as on the previous slide.



General Fluxes

$$\frac{\partial(x)}{\partial t} = -\frac{1}{c_t} \frac{\partial(\mathbf{J})}{\partial z} \quad (\mathbf{J}) = -c_t[D] \left(\frac{\partial x}{\partial z} \right)$$

What have we assumed here?



$$\frac{\partial(x)_i}{\partial t} = -\frac{1}{c_t} \frac{(\mathbf{J})_{i+\frac{1}{2}} - (\mathbf{J})_{i-\frac{1}{2}}}{\Delta z}$$

Applies at $i=1 \dots n$.

NOTE: “ i ” denotes a spatial index, not a species index.

$$(\mathbf{J})_{i-\frac{1}{2}} = -c_t[D]_{i-\frac{1}{2}} \frac{(x)_i - (x)_{i-1}}{\Delta z}$$

At $i=1$ we need to apply BCs

$$(\mathbf{J})_{i+\frac{1}{2}} = -c_t[D]_{i+\frac{1}{2}} \frac{(x)_{i+1} - (x)_i}{\Delta z}$$

At $i=n$ we need to apply BCs

$n_{species}-1$ Dimensional:

(x) - species mole fractions

(\mathbf{J}) - species diffusive fluxes

$[D]$ - Fickian diffusion coefficient matrix

Example: Variable Effective Diffusivity

This represents n_s-1 coupled PDEs for the mole fractions, x_i .

$$\begin{aligned}\frac{\partial x_i}{\partial t} &= -\frac{1}{c_t} \nabla \cdot \mathbf{J}_i \\ &= \nabla \cdot (D_{i,eff} \nabla x_i)\end{aligned}$$

What are the assumptions?

Assume Neumann BCs: $\frac{\partial x_i}{\partial z} = 0$ @ $z=0$ and $z=L$

Given D_{ij} , $x_i(t=0)$, how would you solve this?